

# Judd ofelt analysis of the $\text{Er}^{3+}$ ( $4f^{11}$ ) absorption intensities in $\text{Er}^{3+}$ -doped $\text{Sc}_2\text{O}_3$

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A spectroscopic study on  $\text{Er}^{3+}$  ( $4f^{11}$ ) ions doped in  $\text{Sc}_2\text{O}_3$  transparent ceramics was performed in order to assess its potential as a laser material for different emission wavelengths. The absolute intensities of the absorption bands in the 370 to 1600 nm range were measured at 300 K. The  $\Omega_2, \Omega_4, \Omega_6$ , Judd-Ofelt intensity parameters for  $\text{Er}^{3+}$   $f-f$  transitions were determined and are comparable with the experimental parameters reported for  $\text{Er}^{3+}$  in other sesquioxides single crystals.

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## 1. Introduction

$\text{Er}^{3+}$  doped laser materials are interesting for the development of lasers at different wavelengths from visible to near infrared (especially  $\sim 1.5 \mu\text{m}$  [1,2] and  $\sim 3 \mu\text{m}$  [3]) necessary for a series of applications. The laser emission at  $\sim 1.5 \mu\text{m}$  is used in fiber optics and optical amplifiers, for telecommunication or medicine. The visible emissions in green or red are attractive for upconversion laser devices pumped with infrared laser diodes [4-7]. The cubic sesquioxides, especially scandium oxide ( $\text{Sc}_2\text{O}_3$ ) crystals with their very high thermal conductivity, chemical stability, optical properties (low phonons) and the possibility of doping with  $\text{RE}^{3+}$  are interesting as laser materials. However, due to their very high melting point  $\sim 2430^\circ\text{C}$  [8], it is very difficult to grow such crystals with high optical quality and large sizes. In the recent years the transparent ceramics techniques have been developed to overcome this difficulty. Our investigation deals with the spectral properties of  $\text{Er}^{3+}$  doped  $\text{Sc}_2\text{O}_3$  transparent ceramics.

The  $\text{Sc}_2\text{O}_3$  – scandium sesquioxide forms at room temperature a cubic C-type structure, belonging to the  $I\bar{A}3$  space group. The unit cell contains 16 formula units with 32 cations that form 24 sites of  $C_2$  symmetry and 8 sites of  $C_{3i}$  symmetry. The  $C_2$  site is an eightfold cubic structure with two oxygen vacancies on a face diagonal, while  $C_{3i}$  correspond to a cube with two vacancies on a body diagonal. The lattice constant is  $\sim 10\text{\AA}$  and the cationic density ( $3.334 \times 10^{22}/\text{cm}^3$ ) is rather high compared to other laser crystals. The  $\text{RE}^{3+}$  dopants are assumed to occupy randomly both sites, but the induced electric dipole transitions are symmetry allowed only for  $C_2$  centers.

Some of the spectroscopic properties of  $\text{Er}^{3+}$  doped  $\text{Sc}_2\text{O}_3$  transparent ceramics have been recently investigated in our group [9, 10]. In the present study, additional spectral parameters are determined in order to assess the perspectives of  $\text{Er}^{3+}$  doped  $\text{Sc}_2\text{O}_3$  transparent ceramics as laser material, the Judd-Ofelt [11, 12] (JO) parameters. The Judd-Ofelt parameters are important for the understanding of the relationships between host and  $\text{Er}^{3+}$  emission properties. No JO analysis has been published for Er:  $\text{Sc}_2\text{O}_3$  transparent ceramics, or even for single crystals.

## 2. Experimental methods

The Er-doped  $\text{Sc}_2\text{O}_3$  transparent ceramics were prepared at Japan Fine Ceramics Center, Nagoya, Japan by solid-state synthesis followed by isostatic compression ( $\sim 2$  MPa) and vacuum sintering at  $1750^\circ\text{C}$ . The average size of grains is in the  $30 \mu\text{m}$  range. The optical spectroscopic measurements of 0.3 at% Er:  $\text{Sc}_2\text{O}_3$  were performed on an extended spectral range  $6000\text{cm}^{-1} - 28000 \text{cm}^{-1}$ . The room temperature absorption spectra were measured with a set-up consisting of a Jarell Ash monochromator, S20 and S1 photomultipliers, Si and Ge photodiodes and a Lock in amplifier SR830 on line with a computer.

## 3. Results and discussion

A series of spectral parameters of  $\text{Er}^{3+}$  ions in the  $\text{Sc}_2\text{O}_3$  transparent ceramics could be estimated on the basis of the Judd-Ofelt (JO) theory, which is used largely in evaluating the probability of the forced electric-dipole

transitions of rare-earth ions in various environments, and is often used to calculate the spectroscopic parameters, among them JO parameters. These parameters can be regarded as the phenomenological intensity parameters that characterize the radiative transition probability. The line strengths of the transitions between the ground  $^4I_{15/2}$   $\text{Er}^{3+}$  manifold and the excited  $J'$  manifolds can be calculated. According to the Judd-Ofelt theory, the line strength for forced electric dipole (e.d.) transitions between two manifolds characterized by their total angular momentum quantum numbers,  $J$   $|4f^n[S, L]J\rangle$  and  $J'$   $|4f^n[S', L']J'\rangle$  can be expressed in the following form:

$$S_{calc}^{ed} = \sum_{t=2,4,6} \Omega_t \left\langle 4f^n[S, L]J \left\| U^{(t)} \right\| 4f^n[S', L']J' \right\rangle^2, \quad (1)$$

where  $\Omega_{(t)}$  are JO parameters,  $\left\langle U^{(t)} \right\rangle$  are the doubly reduced matrix elements of unit tensor operator calculated in the intermediate-coupling approximation. The values of the squared reduced elements for the chosen  $\text{Er}^{3+}$  transitions were taken from Carnall *et al.* [13]. The squared matrix elements are independent of the crystal host. The measured line strength,  $S_{meas}(J \rightarrow J')$  of the chosen bands are determined using the following expression:

$$S_{meas}(J \rightarrow J') = \frac{3ch(2J+1)n}{N_0 3\pi^3 \lambda_0} \left[ \frac{9}{(n^2+2)^2} \right] \Gamma, \quad (2)$$

where  $c$  is the speed of light,  $h$  is the Planck constant,  $N_0$  is the  $\text{Er}^{3+}$  ion concentration,  $n$  is the bulk index of refraction,  $\bar{\lambda}$  is the mean wavelength of the absorption band that corresponds to the  $J \rightarrow J'$  transition,  $\Gamma = \int k(\lambda) d(\lambda)$  is the integrated absorption coefficient,  $k(\lambda)$  is the absorption coefficient function on  $\lambda$ . The factor  $[9/(n^2+2)^2]$  represents the local field correction for the ion in a dielectric medium. From several experimental refractive indices of the bulk  $\text{Sc}_2\text{O}_3$  given in [14, 15], the wavelengths dependence of the refractive indices was obtained by using a least-squared fitting program to the Sellmeier dispersion equation:

$$n^2(\lambda) = 1 + \frac{S\lambda^2}{\lambda^2 - \lambda_0^2}, \quad (3)$$

That provides the  $S = 2.856$  and  $\lambda_0 = 76.896$  nm constants, used to recalculate the values of the indices of refraction at the appropriate mean wavelengths of the  $\text{Er}^{3+}$  ion absorption.

To determine the Judd Ofelt (JO) parameters for

$\text{Er}^{3+}$  ( $4f^{11}$ ) transitions in  $\text{Sc}_2\text{O}_3$  transparent ceramic eleven absorption bands of the room temperature spectra were used. In Fig. 1-4 the room temperature absorption for  $C_2$  centers in  $\text{Er}:\text{Sc}_2\text{O}_3$  transparent ceramic ranging from 450 to 1600nm are presented. Only  $C_2$  centers (3/4 from the actual  $\text{Er}^{3+}$  concentration) contribute to e.d. transitions. For  $^4I_{15/2} \rightarrow ^4I_{13/2}$  transition contain the magnetic dipole contribution too.

In order to estimate the JO parameters the experimental line strengths were measured using eq. (2). The values of the integrated absorbance and the mean wavelength were calculated after subtracting the base line of the spectrum. This base line subtraction affects the accuracy of the band integration and thus the accuracy of  $\Omega_t$  parameters determination.

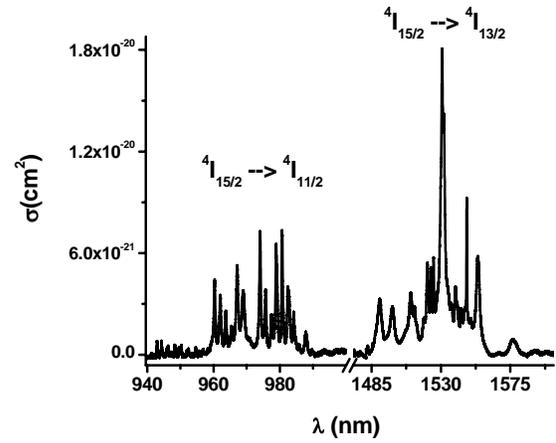


Fig.1  $^4I_{15/2} \rightarrow ^4I_{13/2}$ ,  $^4I_{15/2} \rightarrow ^4I_{11/2}$  absorption cross-sections spectrum of 0.3at. % Er:  $\text{Sc}_2\text{O}_3$  at 300K.

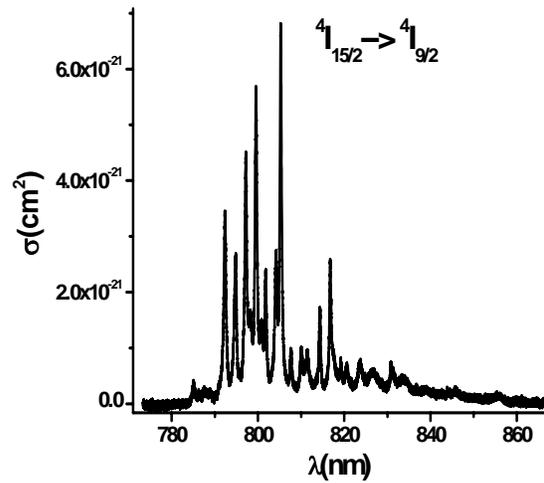


Fig.2  $^4I_{15/2} \rightarrow ^4I_{9/2}$  absorption spectrum of 0.3at. % Er:  $\text{Sc}_2\text{O}_3$  at 300K

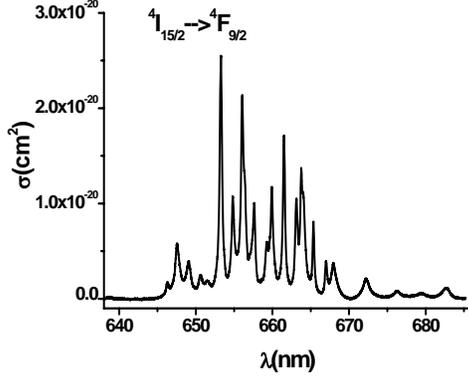


Fig.3  $^4I_{15/2} \rightarrow ^4F_{9/2}$  absorption spectrum of 0.3at. % Er:  $\text{Sc}_2\text{O}_3$  at 300K

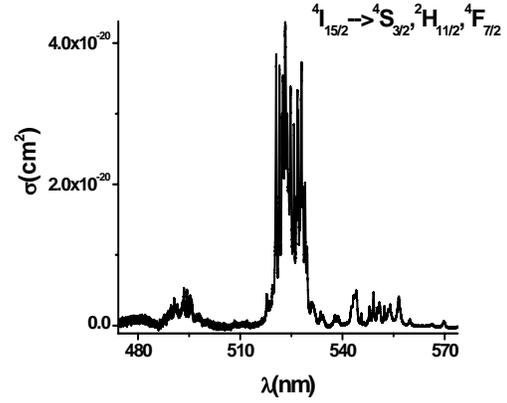


Fig.4  $^4I_{15/2} \rightarrow ^4S_{3/2}, ^2H_{11/2}, ^4F_{7/2}$  absorption spectrum of 0.3at. % Er:  $\text{Sc}_2\text{O}_3$  at 300K

Table I. Measured and calculated absorption line strength of  $\text{Er}^{3+}$  in  $\text{Sc}_2\text{O}_3$  transparent ceramic at 300K.

Transition from $^4I_{15/2}$ to excited states	$\bar{\lambda}$ (nm)	$n$	$\Gamma$ ( $\text{nm cm}^{-1}$ )	$S_{\text{meas}}$ ( $10^{-20} \text{cm}^2$ )	$S_{\text{calc}}$ ( $10^{-20} \text{cm}^2$ )
$^4I_{13/2}$	1512.5	1.965	28.53	1.553	1.398
$^4I_{11/2}$	976.86	1.968	5.105	0.445	0.421
$^4I_{9/2}$	802.9	1.97	3.11	0.326	0.300
$^4F_{9/2}$	658.99	1.973	9.8	1.264	1.216
$^4S_{3/2}$	552.63	1.978	1.54	0.235	0.200
$^2H_{11/2}$	525.84	1.979	20.1	3.236	3.272
$^4F_{7/2}$	496.35	1.981	1.937	0.329	0.742
$^4F_{5/2}, ^4F_{3/2}$	457.76	1.984	1.02	0.188	0.282
$^2G_{9/2}, ^4F_{9/2}$	409.35	1.99	1.34	0.274	0.215
$^4G_{11/2}$	382.97	1.994	19.5	4.26	4.236
$^4G_{9/2}$	369.13	1.996	2.39	0.541	0.487

In Table I the mean wavelength, refractive indices, integrated absorption coefficients and measured and calculated line strength for several  $\text{Er}^{3+}$  transitions in  $\text{Sc}_2\text{O}_3$  ceramic are given.

The intensity JO parameters  $\Omega_2, \Omega_4, \Omega_6$  for  $\text{Er}^{3+}$  in  $\text{Sc}_2\text{O}_3$  were determined by the procedure given in [16] and are given in Table II. As observed, the JO parameters obtained for  $\text{Er}^{3+}$  in  $\text{Sc}_2\text{O}_3$  transparent ceramics are comparable with those obtained for  $\text{Er}^{3+}$ :  $\text{Y}_2\text{O}_3$  single crystals [16].

Table 2. The JO parameters of  $\text{Er}^{3+}$ :  $\text{Sc}_2\text{O}_3$  transparent ceramics (this work) and  $\text{Er}^{3+}$ :  $\text{Y}_2\text{O}_3$  [16].

Host	$\Omega_2$ ( $10^{-20} \text{cm}^2$ )	$\Omega_4$ ( $10^{-20} \text{cm}^2$ )	$\Omega_6$ ( $10^{-20} \text{cm}^2$ )	$X = \Omega_4 / \Omega_6$
$\text{Er}^{3+}$ : $\text{Sc}_2\text{O}_3$	3.57	1.61	0.79	2.03
$\text{Er}^{3+}$ : $\text{Y}_2\text{O}_3$	4.59	1.21	0.48	2.52

The values of the intensity parameters can then be used to recalculate the transition line strengths of the

absorption bands using the Eq. (1). The calculated absorption line strength  $S_{\text{calc}}$  is given in Table I. A measure of the accuracy of the fit is given by rms deviation:

$$\Delta S_{\text{rms}} = \left[ (q - p)^{-1} \sum (\Delta S)^2 \right]^{1/2}, \quad (4)$$

where  $\Delta S = S_{\text{calc}} - S_{\text{meas}}$  is the deviation,  $q$  is the number of the transition analyzed and  $p$  is the number of the parameters, that in this case are 11 and respectively 3. The values in Table I provide an rms deviation of  $0.164 \times 10^{-20} \text{cm}^2$ .

A spectroscopic quality factor  $X = \Omega_4 / \Omega_6$  has been introduced by Kaminskii for  $\text{Nd}^{3+} ^4F_{3/2} - ^4I_1$  transitions [17] as an important characteristic in predicting the stimulated emission for the laser active media. Similarly, one could use the same parameter to evaluate the upconversion emission for  $\text{Er} ^4S_{3/2} - ^4I_{15/2}$  transition in different hosts. According to [18] for  $\text{Er}^{3+}$   $X$  varies from 0.126 to 3.372. In our case,  $\text{Er}^{3+}$ :  $\text{Sc}_2\text{O}_3$  transparent ceramic, this factor is 2.2 (Table II), within this range.

#### 4. Conclusions

Judd-Ofelt analysis of  $\text{Er}^{3+}$  ion optical intensities in  $\text{Sc}_2\text{O}_3$  transparent ceramic was performed, and based on the absolute absorption data the three phenomenological parameters  $\Omega_2, \Omega_4, \Omega_6$  were determined with. The estimated parameters are similar to those of  $\text{Er}^{3+}$  in  $\text{Y}_2\text{O}_3$  single crystals. Similarly to  $\text{Nd}^{3+}$ , a spectroscopic quality factor  $X = \Omega_4 / \Omega_6$  for the  $\text{Er}^{3+}$  green upconversion emission was determined to be 2.2, within the range for other hosts. These parameters can be used to estimate the radiative probabilities, radiative lifetimes, branching ratios for different transitions; that along with experimental lifetimes allow the determination of the quantum efficiencies. Such investigations are in progress.

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